

## Pre-Semester 2010 - Physics Course - Extra Tutorial

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## 1. Spring

A point mass attached to a (massless) spring with force constant  $D$  is initially displaced by  $x_0$  from its equilibrium position (the spring is, say, stretched). It is released at time  $t = 0$ .

- State the equation of motion (for  $x(t)$ ) and solve it.
- At which time  $t_1$  does the oscillator (the spring with mass point) return to its initial state the first time? How would you name such a time?
- What is the velocity  $v(t)$  at time  $t$ ?
- What are the kinetic and potential energies  $E_{\text{kin}}(t)$  and  $E_{\text{pot}}(t)$  at time  $t$ ? Compute their sum  $E_{\text{tot}}$ .

**Solution:**

- Since  $F = -Dx$  the equation of motion reads  $m\ddot{x} + Dx = 0$ . Its solution is  $x(t) = x_0 \cos(\omega t)$  with  $\omega = \sqrt{D/m}$ .
- The mass point reaches its original position when  $\omega t_1 = 2\pi \Rightarrow t_1 = 2\pi/\omega$ . Of course, this is the period of oscillation.
- $v(t) = \dot{x}(t) = -x_0\omega \sin(\omega t)$
- The kinetic energy is  $E_{\text{kin}}(t) = \frac{1}{2}mv(t)^2 = \frac{1}{2}m\omega^2 x_0^2 \sin^2(\omega t) = \frac{1}{2}Dx_0^2 \sin^2(\omega t)$ . The potential energy is  $E_{\text{pot}}(t) = \frac{1}{2}Dx(t)^2 = \frac{1}{2}Dx_0^2 \cos^2(\omega t)$ . While both oscillate, their sum  $E_{\text{tot}}(t) = E_{\text{kin}}(t) + E_{\text{pot}}(t) = \frac{1}{2}Dx_0^2$  is constant (using  $\sin^2 x + \cos^2 x = 1$ ). It equals the *initial potential energy*.

## 2. Weakly Damped Harmonic Oscillator

We now assume that the oscillator from the previous exercise moves in a liquid and experiences Stokes friction with coefficient  $\gamma$ . We assume that this damping is weak,  $\gamma^2 < 4Dm$ .

- State the new equation of motion (again for  $x(t)$ ).
- Consider the function

$$x(t) = Ae^{-\lambda t} \cos(\omega t + \phi) \quad \text{with} \quad \lambda = \frac{\gamma}{2m} \quad \text{and} \quad \omega = \sqrt{\frac{D}{m} - \frac{\gamma^2}{4m^2}}.$$

Compute its derivatives  $\dot{x}(t)$ ,  $\ddot{x}(t)$  and show that it solves the equation of motion.

**Solution:**

(a) The friction force is  $F_s = -\gamma v = -\gamma \dot{x}$  which yields the equation of motion

$$m\ddot{x} + \gamma\dot{x} + Dx = 0.$$

(b) Function and derivatives are

$$\begin{aligned} x(t) &= Ae^{-\lambda t} \cos(\omega t + \phi), \\ \dot{x}(t) &= Ae^{-\lambda t} [-\lambda \cos(\omega t + \phi) - \omega \sin(\omega t + \phi)], \\ \ddot{x}(t) &= Ae^{-\lambda t} [(\lambda^2 - \omega^2) \cos(\omega t + \phi) + 2\lambda\omega \sin(\omega t + \phi)]. \end{aligned}$$

Substituting them into the equation of motion and collecting terms with cos and sin gives

$$0 = Ae^{\lambda t} \{ [m(\lambda^2 - \omega^2) - \lambda\gamma + D] \cos(\omega t + \phi) + [2m\lambda\omega - \gamma\omega] \sin(\omega t + \phi) \}.$$

The square bracket in front of  $\sin(\dots)$  vanishes if

$$\lambda = \frac{\gamma}{2m}.$$

Then the square bracket in front of  $\cos(\dots)$  vanishes if

$$\lambda^2 - \omega^2 = \frac{\lambda\gamma - D}{m} \Leftrightarrow \omega = \pm \sqrt{\lambda^2 - \frac{\lambda\gamma - D}{m}} = \pm \sqrt{\frac{D}{m} - \frac{\gamma^2}{4m^2}}.$$

**3. Physical Pendulum**

At time  $t = 0$  a physical pendulum with moment of inertia  $\Theta = 0.22 \text{ kg m}^2$  and mass  $m = 500 \text{ g}$  makes an angle  $\alpha_0 = 7^\circ$  (more precisely: the line connecting rotation axis and center of mass makes this angle) with the vertical and has an angular velocity  $\omega_0 > 5.75^\circ/\text{s}$  (it is moving counterclockwise). Its rotation axis and center of mass have a distance  $l = 10 \text{ cm}$ .

- State the equation of motion (for  $\alpha(t)$ ) and approximate it for small angles  $\alpha$ .
- Solve the equation of motion by using the general ansatz and the *two* initial conditions.
- What is the amplitude of the oscillation (in degrees)? What is the period of oscillation?
- At which time  $t_1$  will the pendulum reach maximal deflection (maximal  $\alpha$ ) for the first time? At which time  $t_2$  will it do so for the second time (this time deflection will be negative)?
- Assuming that the pendulum has been oscillating for quite some while (before  $t = 0$ ), at which time  $t_0$  was the angle  $\alpha = 0$  the last time before  $t = 0$ ?
- Sketch the solution.

**Solution:**

(a) With our sign convention moment of torque and angular momentum read

$$M(t) = -mgl \sin \alpha(t), \quad L(t) = \Theta \omega(t) = \Theta \dot{\alpha}(t).$$

Using the angular version of Newton's second law,  $L = \dot{M}$ , we obtain for the equation of motion:

$$\Theta \ddot{\alpha}(t) + mgl \sin(\alpha(t)) = 0.$$

Assuming that the angle  $\alpha$  is small (for all times  $t$ ), we may approximate the equation:

$$\ddot{\alpha}(t) + \frac{mgl}{\Theta} \alpha(t) = 0.$$

(b) Expecting harmonic oscillation, we use the ansatz

$$\alpha(t) = A \cos(\omega t + \phi).$$

Then,

$$\omega = \sqrt{\frac{mgl}{\Theta}} = 1.5/\text{s} \quad \text{and} \quad \dot{\alpha}(t) = -A\omega \sin(\omega t + \phi).$$

The two unknown quantities  $A$  and  $\phi$  are determined by the two initial conditions:

$$\alpha(0) = A \cos(\phi) = \alpha_0, \tag{1}$$

$$\dot{\alpha}(0) = -A\omega \sin(\phi) = \omega_0. \tag{2}$$

$$(2)/(1) \Rightarrow -\omega \tan(\phi) = \frac{\omega_0}{\alpha_0} \Rightarrow \phi = -\arctan \left[ \frac{\omega_0}{\alpha_0} \frac{1}{\omega} \right] = -0.5,$$

$$(1) \Rightarrow A = \frac{\alpha_0}{\cos(\phi)} = 8^\circ,$$

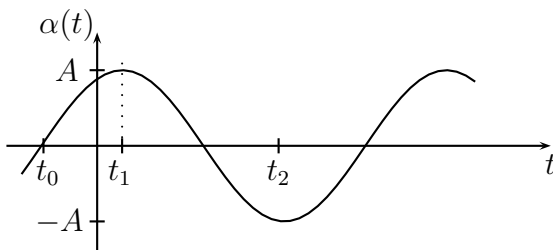
$$\Rightarrow \boxed{\alpha(t) = 8^\circ \cos(1.5t/\text{s} - 0.5)}.$$

(c)  $A = 8^\circ$ ,  $T = \frac{2\pi}{\omega} = 4.2\text{s}$ .

(d)  $\boxed{t_1}$   $\omega t_1 + \phi = 0 \Rightarrow t_1 = -\phi/\omega = 0.33 \text{ s}$ .

$\boxed{t_2}$   $t_2 = t_1 + T/2 = 2.4 \text{ s}$ .

(e)  $\boxed{t_0}$   $\omega t_0 + \phi = -\frac{\pi}{2} \Rightarrow t_0 = (-\frac{\pi}{2} - \phi)/\omega = -0.71 \text{ s}$ .



(f)

**4. Where's the Angle?**

A point mass  $m$  is performing a circular motion in the two dimensional  $x - y$ -plane around the origin  $O$ . It moves counterclockwise with constant angular velocity  $\omega$  along a circle with radius  $A$ .

- (a) What is the angle  $\alpha(t)$  at time  $t$  between the line connecting  $O$  with the point mass and the  $x$ -axis if for  $t = 0$  this angle is  $\phi$ ?
- (b) What are the coordinates  $x(t)$  and  $y(t)$  of the point mass at time  $t$ ?
- (c) Do you remember how the point mass's speed  $v$  is related to the given quantities?
- (d) Just to be sure, check that  $v^2 = (\dot{x}(t))^2 + (\dot{y}(t))^2$ .
- (e) Give the force  $F_{\text{cp}}$  which acts on the point mass. What is the component  $F_{\text{cp},x}$  of this force in  $x$ -direction? Express it in terms of  $x$  (get rid of  $\alpha$  or  $t$ !).
- (f) Although it may seem pointless, since we already know the solution: Using Newton's second law  $m\ddot{x} = F_{\text{cp},x}$  in  $x$ -direction, state the equation of motion for  $x(t)$ .

**Solution:**

- (a)  $\alpha(t) = \omega t + \phi$
- (b)  $x(t) = A \cos(\alpha(t)) = A \cos(\omega t + \phi)$ ,  $y(t) = A \sin(\alpha(t)) = A \sin(\omega t + \phi)$ .
- (c)  $v = \omega A$
- (d) It is  $\dot{x}(t) = -A\omega \sin(\omega t + \phi)$  and  $\dot{y}(t) = A\omega \cos(\omega t + \phi)$ . Since  $\sin(x)^2 + \cos(x)^2 = 1$  it is

$$(\dot{x}(t))^2 + (\dot{y}(t))^2 = A^2\omega^2 \sin^2(\omega t + \phi) + A^2\omega^2 \cos^2(\omega t + \phi) = A^2\omega^2 = v^2.$$

- (e) The centripetal force is  $F_{\text{cp}} = mv^2/A = mA\omega^2$ . Its  $x$ -component is

$$F_{\text{cp},x} = -\cos(\alpha) F_{\text{cp}} = -A \cos(\alpha) m\omega^2 = -m\omega^2 x.$$

The sign ensures that if  $x$  is positive, the force points to the left, hence is negative.

- (f) The equation of motion thus is  $m\ddot{x} + \omega^2 x = 0$ .

In the considered situation we have a circular motion and the notions of an angle  $\alpha = \omega t + \phi$  and of angular velocity completely make sense. However, we have seen that if we consider only the  $x$ -coordinate, that means, if we follow the "shadow" of the point mass on the  $x$ -axis, we obtain a harmonic oscillation.

It is this analogy which justifies the notion of *angular velocity*  $\omega$  for all harmonic oscillations even if there is no actual angle  $\alpha$  involved. Usually, one then refers to  $\alpha$  as the *phase*. As the angle  $\alpha(t)$  marked the position of our point mass at a given time  $t$ , the more abstract phase  $\alpha(t)$  characterizes the state of any harmonic oscillator at time  $t$ .